



*FI<sup>3</sup>*

# Holistic Simulation for FIREX Project with FI<sup>3</sup>

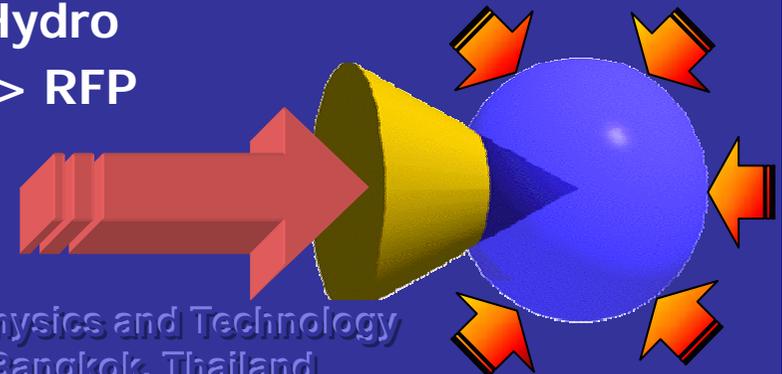
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## Outline

1. Introduction (FI<sup>3</sup>, previous works)
2. Fast Electron Generation ... PIC
  - a. Effect of Density gap at rear surface of cone tip
  - b. Density profile steepening at LPI region
  - c. Pre-plasma scale length dependence
3. Core Heating .....RFP-Hydro
4. Fast Ion Contribution ..... PIC -> RFP
5. Summary

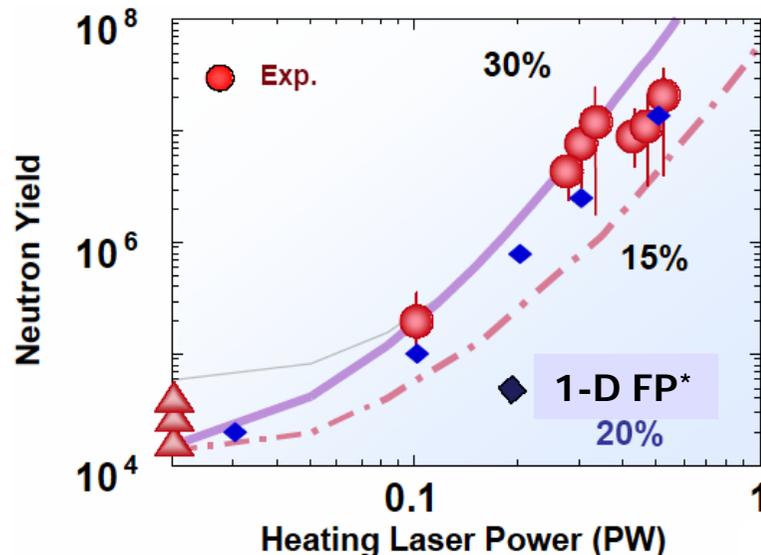


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March 5-9, 2007, Amari Atrium Hotel, Bangkok, Thailand

# Fast Heating Experiments at ILE

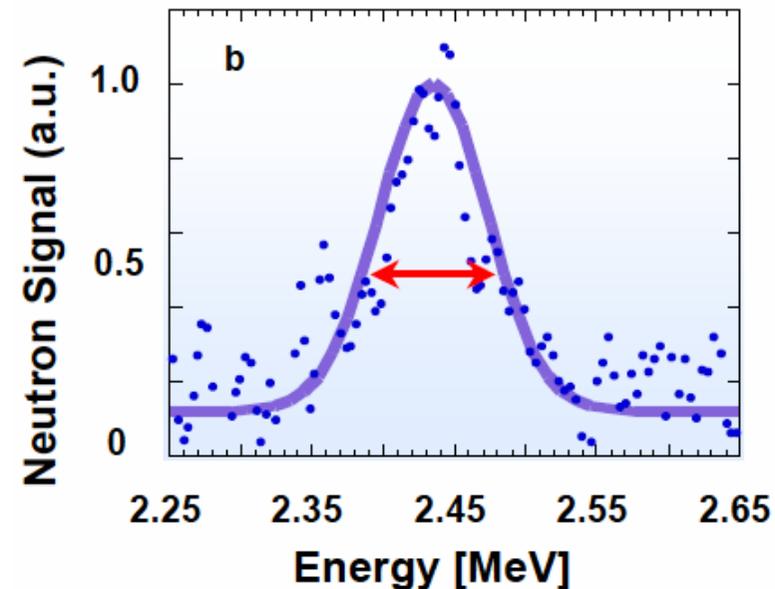
## Fast Heating of Cone-Guided CD Targets with GEKKO PW Laser at ILE R. Kodama, et al., Nature 418, 933 (2002)

Neutron yield was enhanced by 3 orders of magnitude.



\*) 1-D Fokker-Planck simulations results

Heated core up to ~ 0.8keV.



As the next step, FIREX project has been started.  
However, efficient heating mechanism is not clarified yet!!

# FI<sup>3</sup> Project

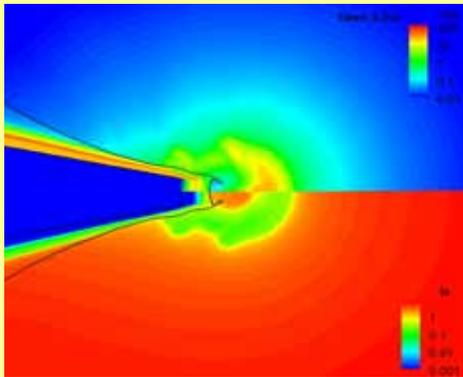
FI<sup>3</sup>

Fast Ignition Integrated Interconnecting code Project

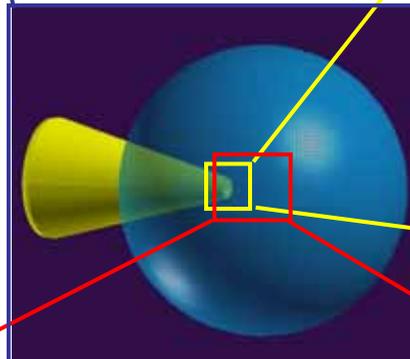
H. Sakagami and K. Mima, Laser Part. Beams, 22 (2004) 41.

Radiation-Hydro code "PINOCO"

Implosion Dynamics

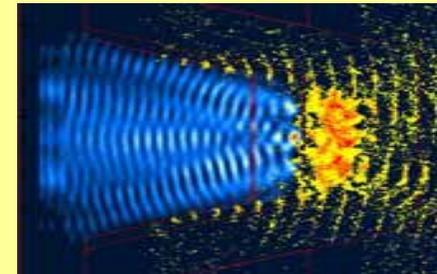


cone - core profiles



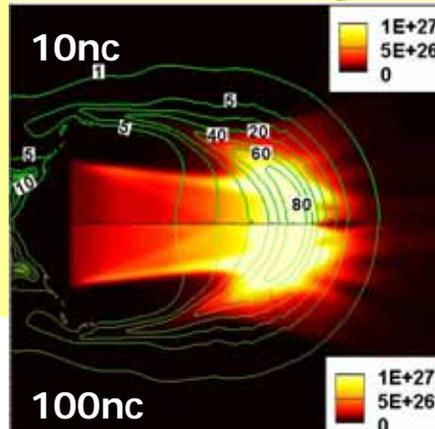
Collective PIC code "FISCOF"

Relativistic LPI



Fokker-Planck - Hydro code "FIBMET"

Core heating



Bulk Plasma profiles  
Energy deposition

Return current  
Field profiles  
Fast electron profiles  
Field profiles

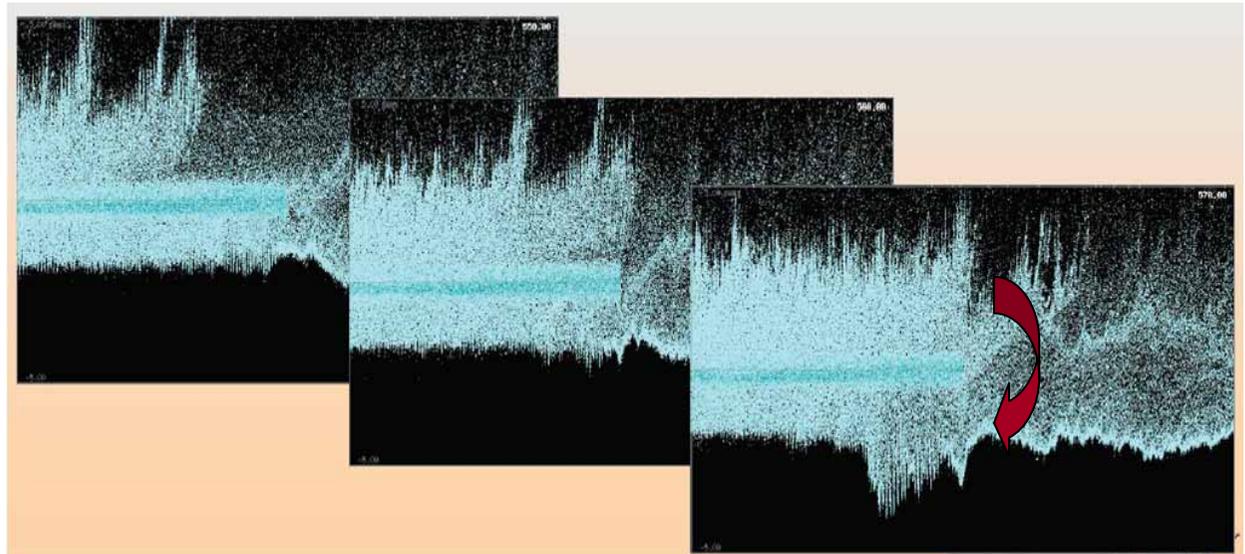
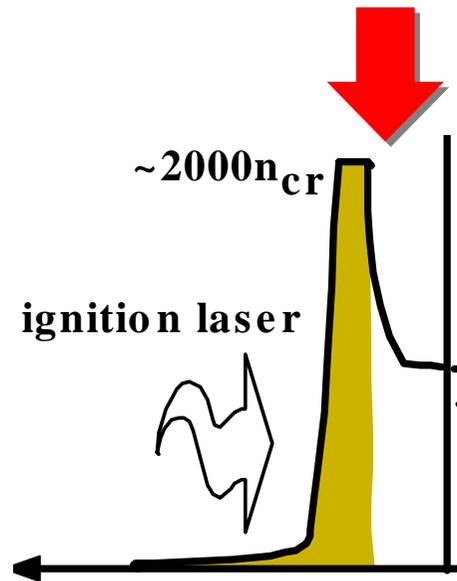
Osaka Univ.,  
NIFS,  
Kyushu Univ.  
Setsunan Univ.

Data flow

→ considered  
← planned

# Previous work

## Density gap effect at rear surface of cone tip



Fast electron current generates the strong static field, which traps relatively low energy fast electrons inside the cone tip and accelerates the bulk electrons into the cone tip.

=> Fast electron energy is moderated and the beam duration becomes long.

Due to the density gap effects, the dense core is heated more efficiently. But, the heated core temperature is  $\sim 0.5\text{keV}$ , which is still lower than the experimentally measured value,  $0.8\text{keV}$ .

# Density gap + Pre-plasma scale length

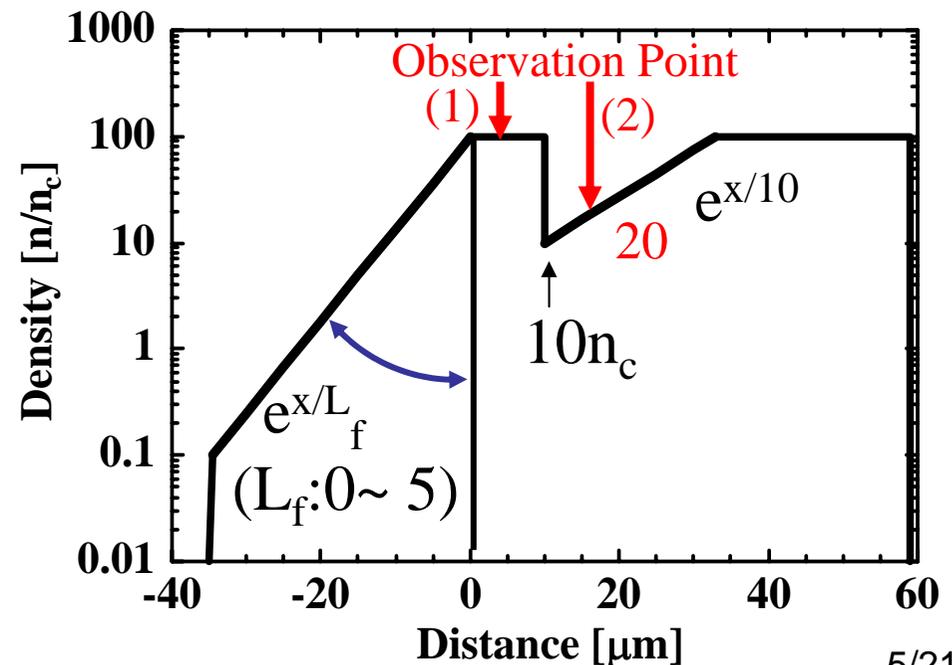
- The pre-plasma scale length, which is determined by the pre-pulse of heating laser, may strongly affect the laser-plasma interactions. (previously fixed as  $5\mu\text{m}$ )
- The effects of density gap + pre-plasma scale length are evaluated by 1D simulations (PIC + RFP-Hydro).

## 1D PIC simulations “FISCOF1”

evaluation of effects of density gap and pre-plasma scale length on the fast electron profiles

## 1D FP-Hydro simulations “FIBMET”

Evaluation of core heating using fast electron profiles obtained by 1D-PIC and bulk plasma profiles obtained by 2-D implosion simulations “PINOCO”

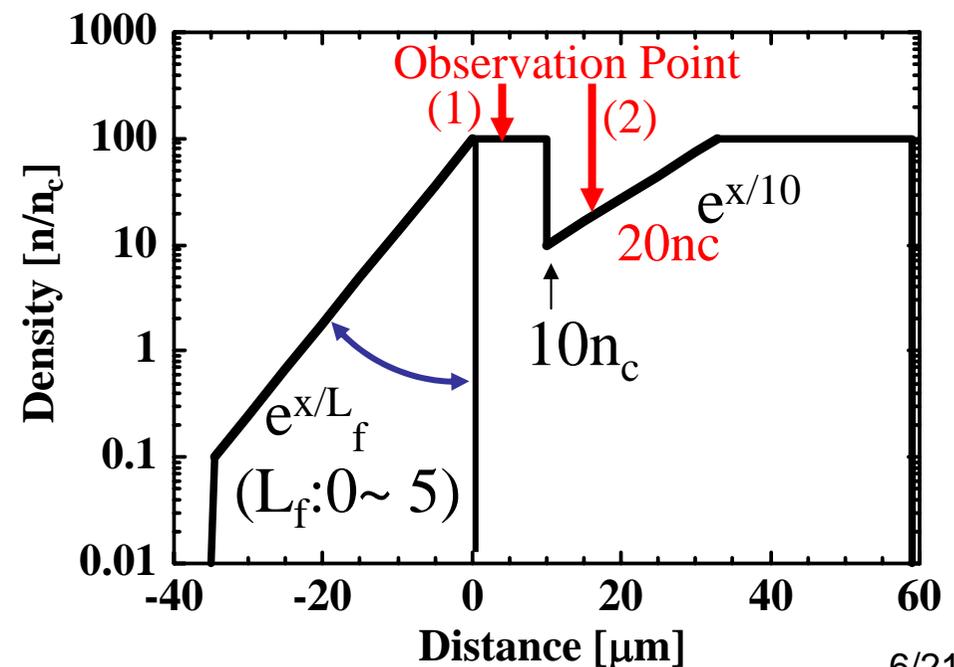


# Simulation Condition for 1D "PIC FISCOF1"

## Laser-Plasma Interactions at the Cone Inner surface

- Electrons and ions are mobile.
- All the ion are Au with real mass and  $Z=50$
- Collision is not included.
- Cone tip is modeled by 100nc plasma
- Heating laser parameters:

- $\lambda_L = 1.06$  [ $\mu\text{m}$ ]
- Gaussian profile
- $\tau_{\text{FWHM}} = 750$  [fs]
- $I_L = 10^{20}$  [ $\text{W}/\text{cm}^2$ ]



# Density Steepening at LPI region

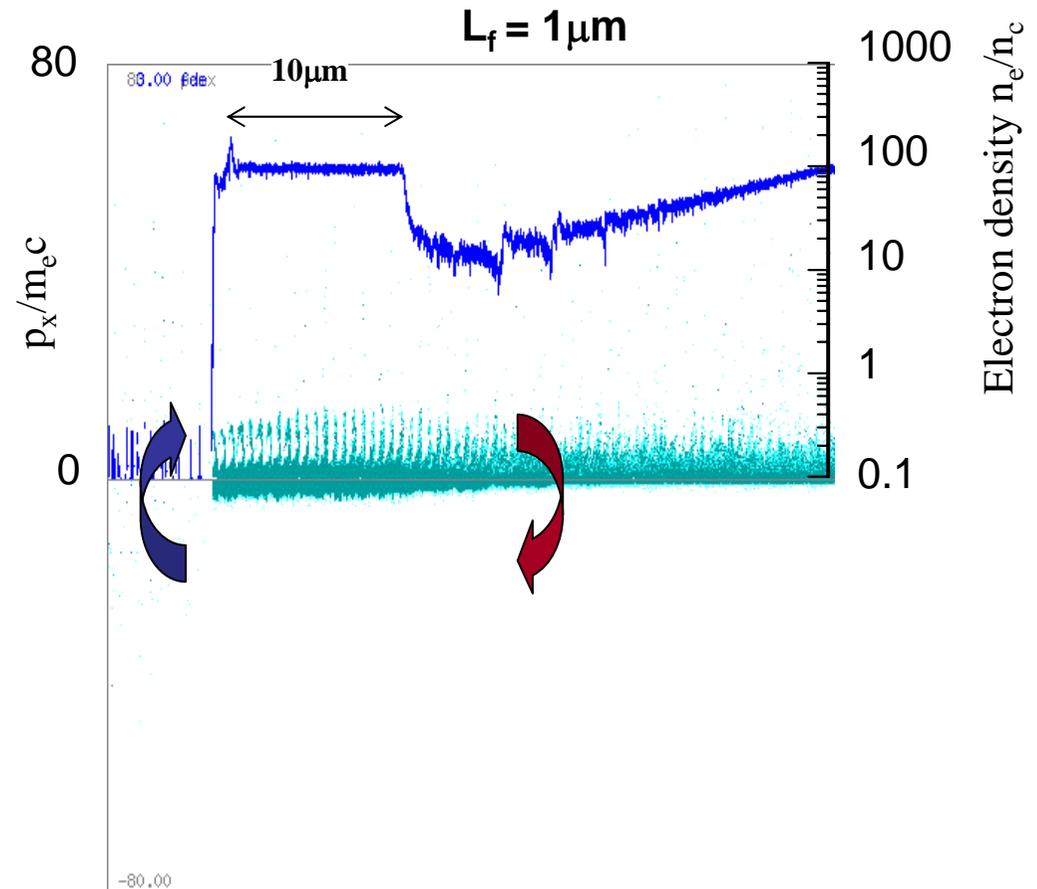
Strong Ponderomotive force steepens the density profile.

## At the front surface (LPI region)

- The pre-plasma is pushed by the strong Ponderomotive force and then the density profile steepening occurs.
- Periodically lunched electron bunches are observed, and returned electrons are completely kicked back.

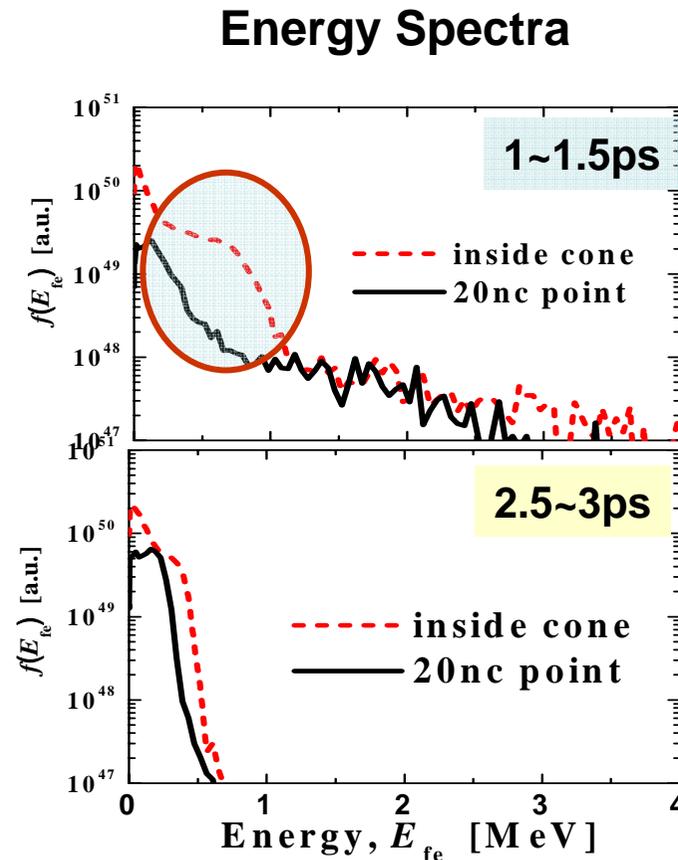
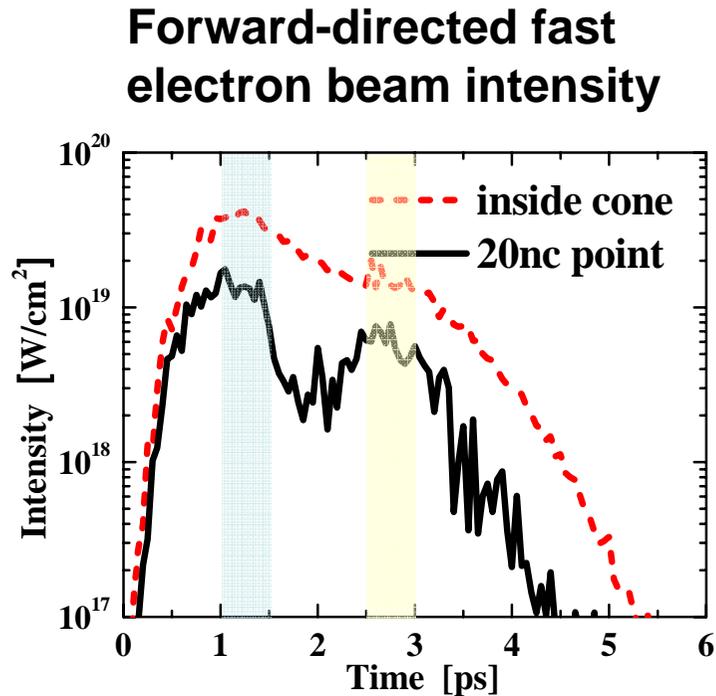
## At the rear surface (density gap)

- The static field built up due to fast electron current, which reflects the relatively low-energy fast electrons and accelerates the bulk electrons into the cone tip.



at  $t = 1\text{ps}$

# Effects of density gap & profile steepening ( $L_f = 1\mu\text{m}$ )



- The relatively low energy fast electrons are trapped inside the cone tip due to the static fields generated at the front and rear surfaces of cone tip.
- These trapped electrons are released from the tip even after finishing laser irradiation.

# Scale length dependence 1

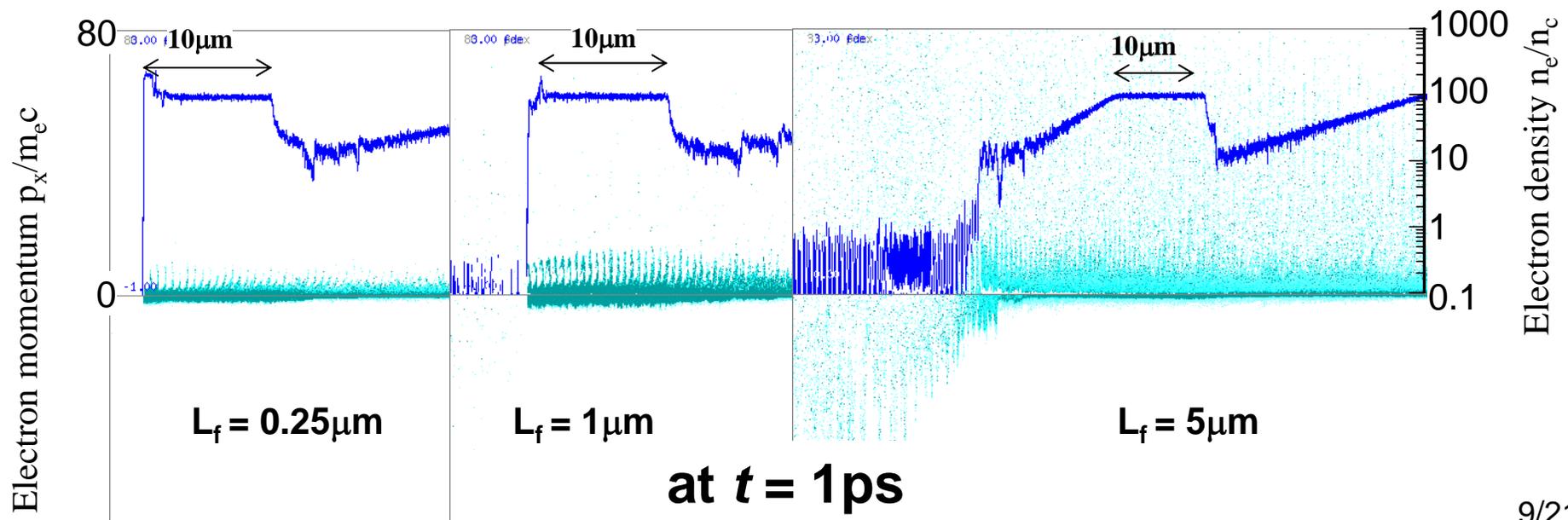
Strong Ponderomotive force steepens the density profile.

## Short scale length case

- The pre-plasma is swapped away.
- Periodically lunched electron bunches are observed, and returned electrons are completely kicked back.

## Long scale length case

- The pre-plasma still remains at LPI region.
- Very high energy fast electrons are generated and part of the returned electrons are pulled into the vacuum region.

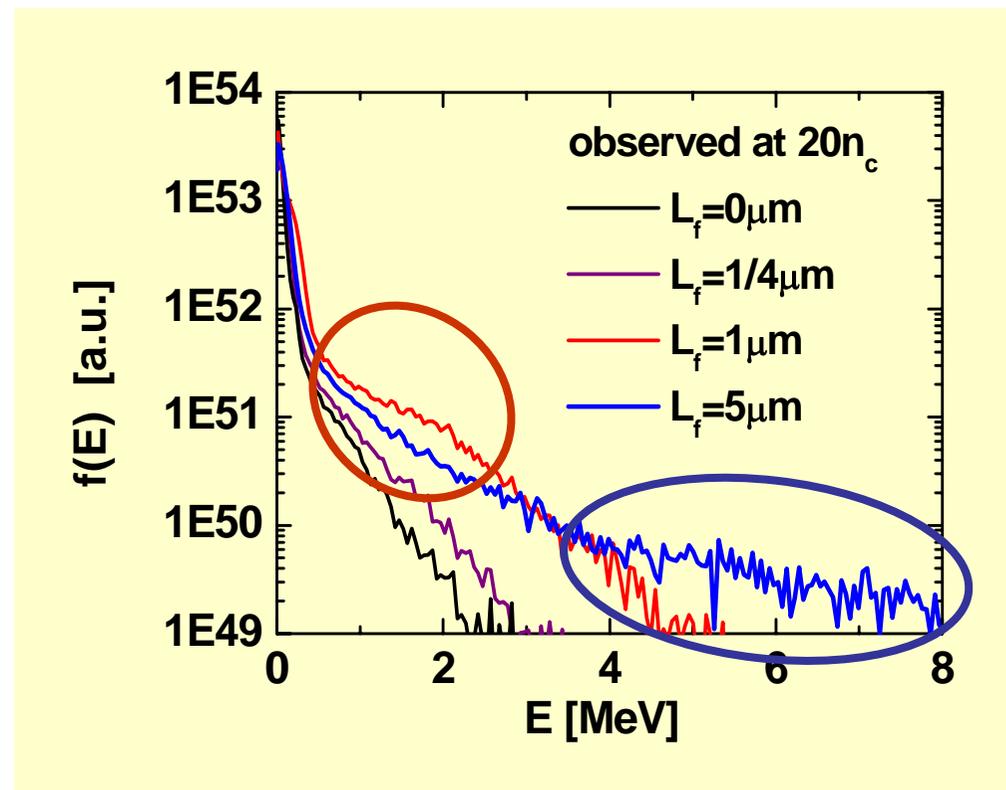


# Scale length dependence 2

With increasing the scale length, higher energy electrons generated.

- $L_f = 5\mu\text{m}$ ; very high energy electrons are not trapped and trapped effect due to density gap and profile steepening is weaker. So the fraction of high energy electrons is large. They can not heat the core efficiently due to long range.
- $L_f = 1\mu\text{m}$ ; the relatively low energy electrons, which are favorable to core heating, are increased due to the gap and steepening effects.

Time-integrated Fast electron spectrum



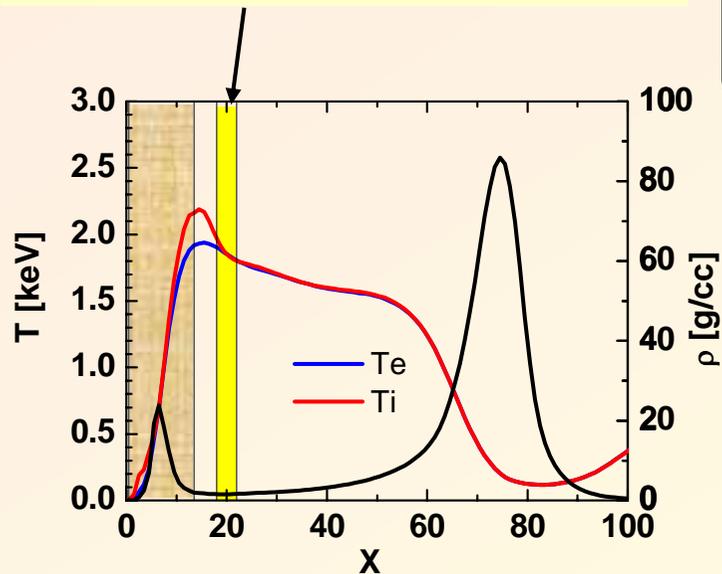
**Optimum scale length for core heating will exist!!**

# 1D Fokker-Planck Simulation Model

Fast Electron & Ion Profiles  
(PIC Code)



*Fast electrons were injected at outer surface of a gold cone.*



Imploded Core Profiles  
(ALE Rad-Hydro Code)



## Fast Electron & Ion Transport

Relativistic Fokker-Planck transport  
Electric Fields

(x) – CIP (85x161 mesh)

(p) – Discontinuous Linear FEM (25 groups)

( $\mu$ ) – Discrete Ordinate Sn method (36 directions)

$n_i, T_i, T_e, Z$

Energy(fast electron, ion )  
& Momentum(fast ion)  
deposition rates

## Radiation-Hydrodynamics

Bulk Plasma

· 1-fluid 2-temp. CIP code

Radiation

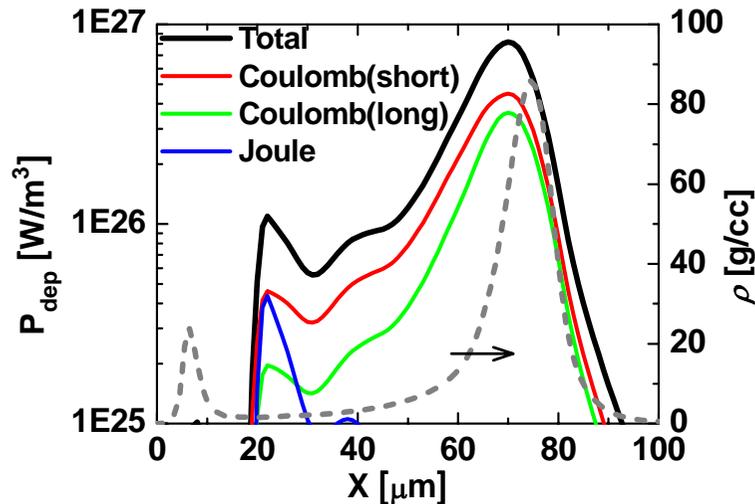
· Flux-limited diffusion

# Core Heating Properties (1D RFP) $F1^3$

## Core Heating Properties...1D FP

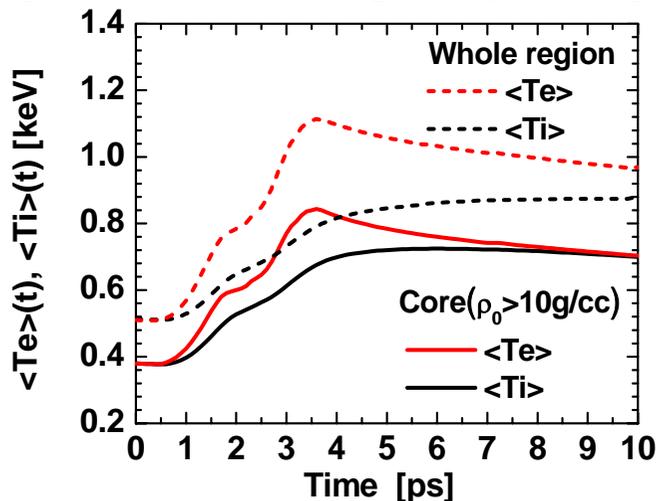
source:  $L_f = 1\mu\text{m}$ ,  $n_{\text{erear}} = 10n_c$

Heating rate (at  $t = 2.5\text{ps}$ )



- The Coulomb interaction with bulk electron is dominant energy deposition mechanism of fast electron in dense core.
- The bulk ion is heated through the temperature relaxation with bulk electron, so Ti reaches maximum by 3ps delay from Te,max time.

Temporal evolution of spatially-averaged Temp



Averaged ion temperature reaches

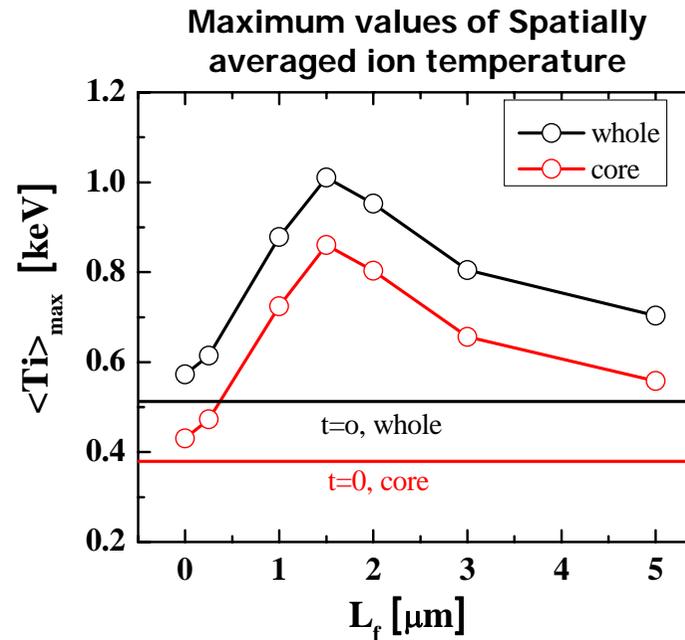
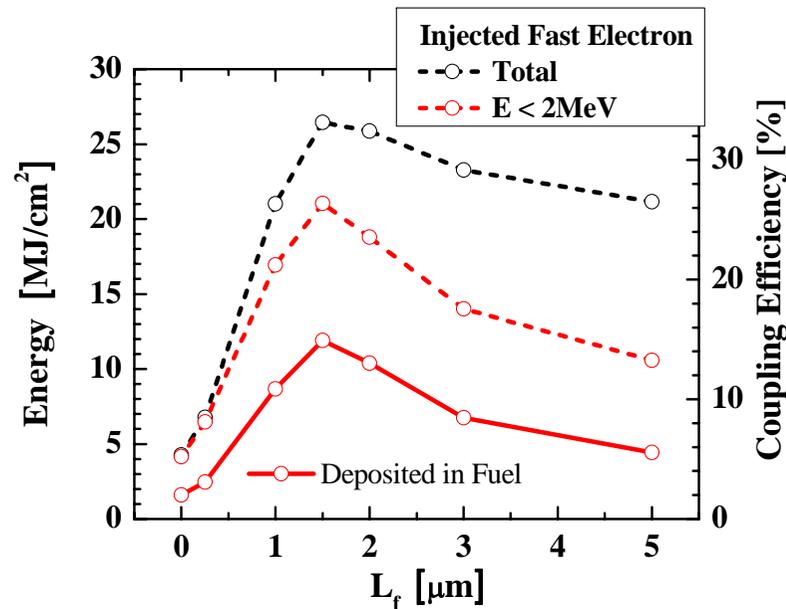
- 0.87keV at 9,0ps in the whole fuel region,
- 0.72keV at 6.5ps in dense core ( $\rho > 10\text{g/cc}$ ).

$$\Delta T_i \sim 0.35\text{keV}$$

# Pre-plasma scale length dependence : $L_f=0 \sim 5 \mu\text{m}$

## Core Heating Properties...1D FP

Irradiated Laser Energy =  $79.8 \text{ MJ/cm}^2$



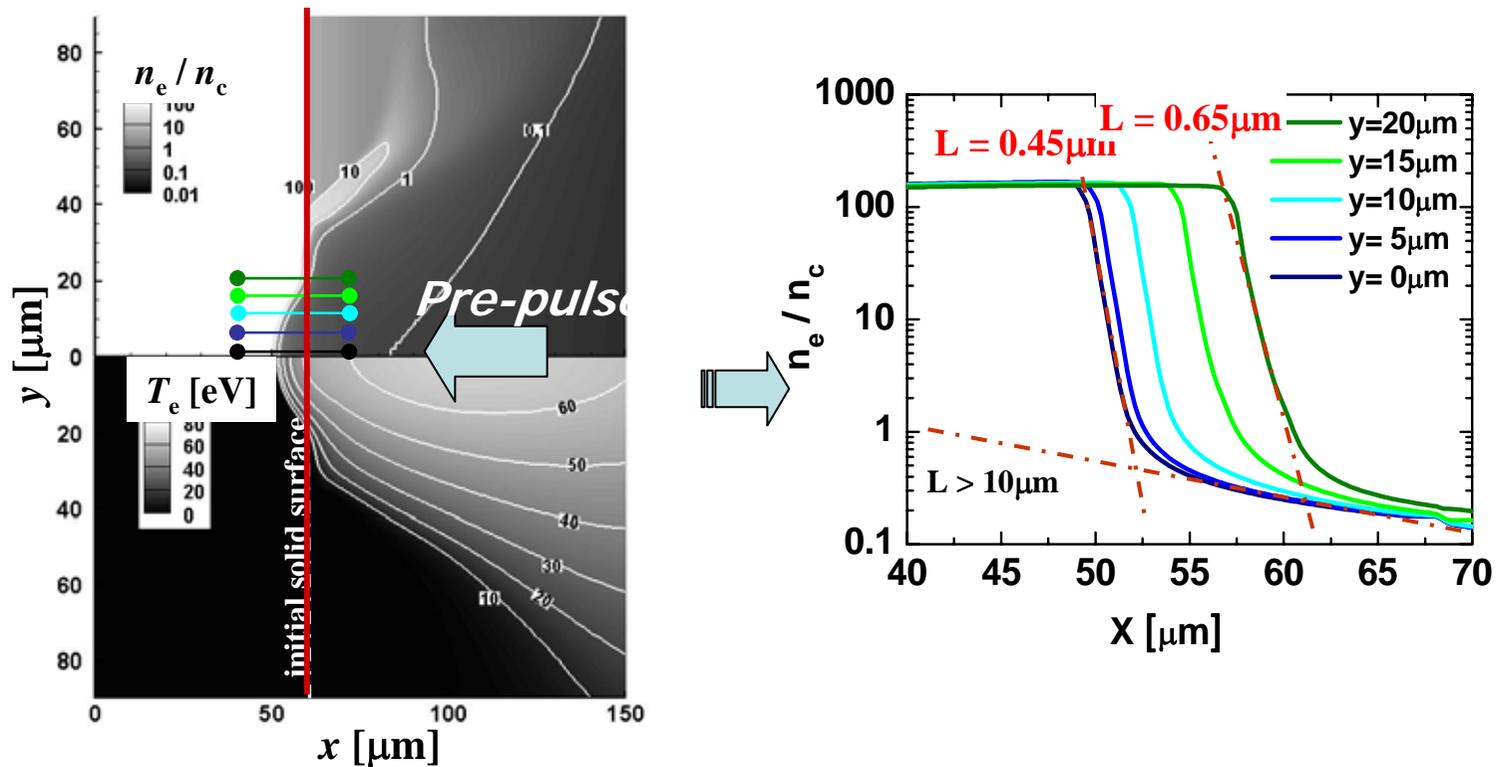
- $L_f < 1 \mu\text{m}$ ;  $\eta_{L_e}$  is small, so that the core heating efficiency is low.
- $L_f > 2 \mu\text{m}$ ; with increasing  $L_f$ ,  $\eta_{L_e}$  gradually decreases, but the low energy component ( $E < 2\text{MeV}$ ) decreases faster, which lowers the core heating efficiency.

In the present simulations, the optimum  $L_f$  for core heating is  $1.5 \mu\text{m}$ . The energy coupling from laser to core is 14.9%, and the core temperature reaches  $0.86 \text{ keV}$  ( $\Delta T_i = 0.48 \text{ keV}$ ).

# Pre-plasma estimation by PINOCO



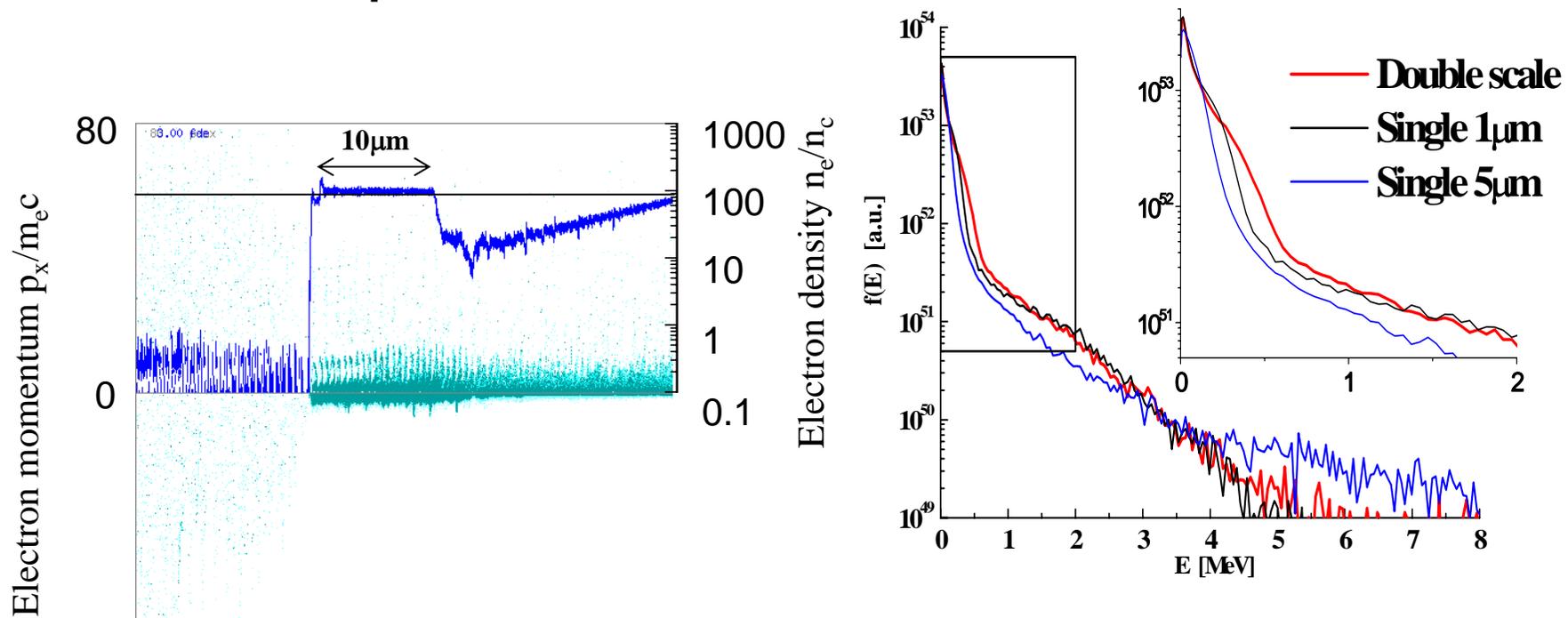
An Example of PINOCO Simulation for pre-plasma formation due to pre-pulse irradiation ( $\sim 10^{11} \text{W/cm}^2$ )



Pre-plasma has double scale length, *i.e.*  $L_{f1} \leq 1 \mu\text{m}$  at  $n_e > n_c$  and  $L_{f2} \gg 1 \mu\text{m}$  at  $n_e < n_c$ .

# Double Scale Length; $(L_{f1} = 1\mu\text{m} (n_e > n_c) + L_{f2} = 5\mu\text{m} (n_e < n_c))$

## Fast electron profiles ..... 1D PIC



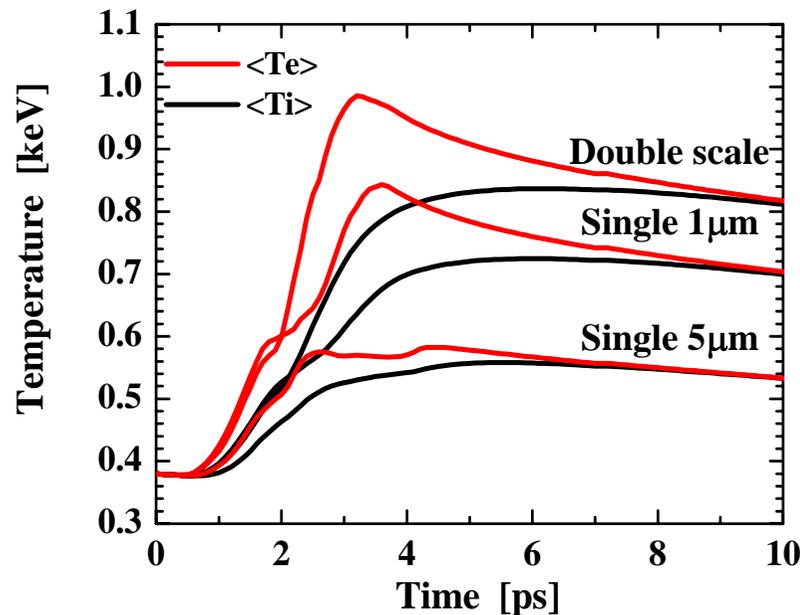
|                      | Electron beam                          | Low energy (2MeV) Comp. |
|----------------------|--|-------------------------|
| $L_f = 1\mu\text{m}$ | 21.0 MJ/cm <sup>2</sup> (26.3%*)       | 16.9 (21.1%*)           |
| $L_f = 5\mu\text{m}$ | 21.2 MJ/cm <sup>2</sup> (26.6%*)       | 10.6 (13.3%*)           |
| <b>Double scale</b>  | <b>24.6 MJ/cm<sup>2</sup> (30.8%*)</b> | <b>19.8 (24.8%*)</b>    |

\* Energy carried by fast particles / laser energy (79.8MJ/cm<sup>2</sup>) [%]

# Double Scale Length; $(L_{f1}=1\mu\text{m} (n_e > n_c) + L_{f2}=5\mu\text{m} (n_e < n_c))$



## Core Heating Properties...1D FP



Irradiated Laser Energy = 79.8 MJ/cm<sup>2</sup>

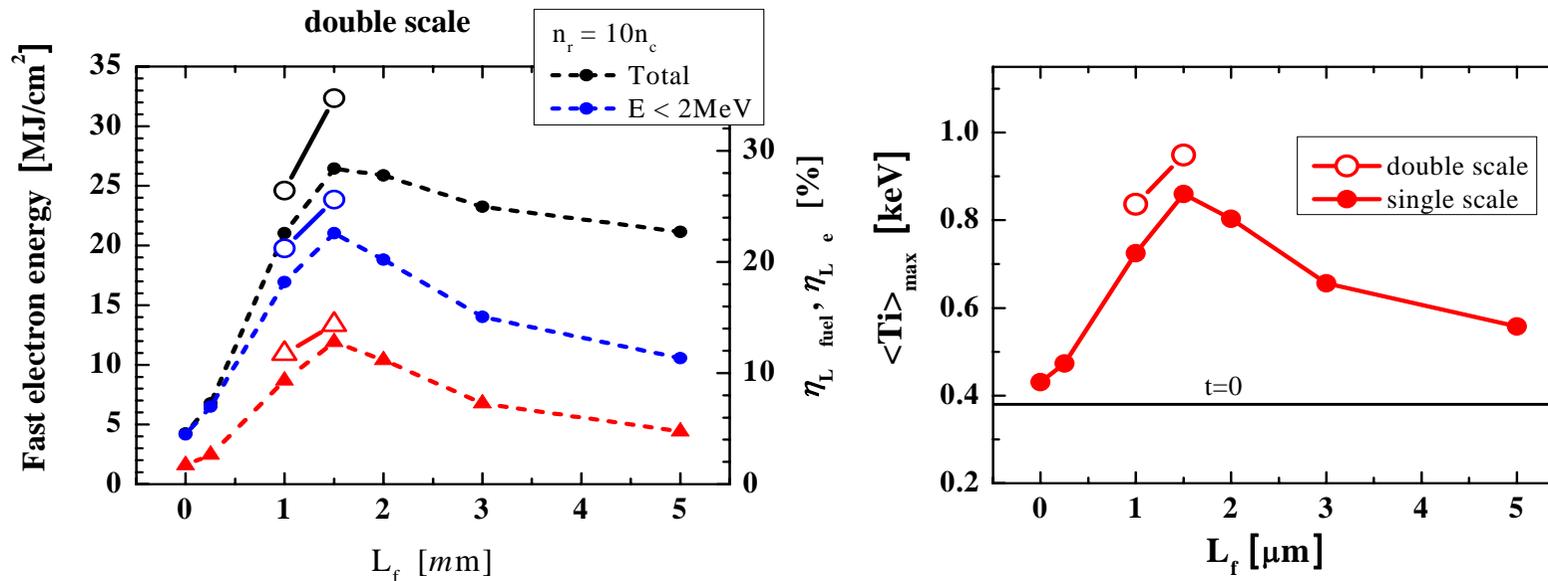
| $L_f[\mu\text{m}]$                                     | <i>Single scale</i> |      | <i>Double scale</i> |
|--|---------------------|------|---------------------|
|  | 1.0                 | 5.0  | 1.0+5.0             |
| $E_{\text{beam}}[\text{MJ}/\text{m}^2]$                | 21.0                | 21.2 | <b>24.6</b>         |
| $\eta_{L_e} [\%]$                                      | 26.3                | 26.6 | <b>30.8</b>         |
| $E_{\text{dep}} [\text{MJ}/\text{m}^2]$                | 8.68                | 4.42 | <b>10.94</b>        |
| $\eta_{e \text{ fuel}} [\%]$                           | 41.3                | 20.8 | <b>44.5</b>         |
| $\eta_{L \text{ fuel}} [\%]$                           | 10.9                | 5.5  | <b>13.7</b>         |
| $\Delta \langle T_i \rangle_{\text{max}} [\text{keV}]$ | 0.34                | 0.17 | <b>0.46</b>         |

In the double scale case, the number of relatively low energy electrons ( $E < 2\text{MeV}$ ) is increased, so that the heated core temperature becomes higher ( $\langle T_i \rangle_{\text{max}} = 0.84\text{keV}$ ) than that in 1μm single scale case ( $\langle T_i \rangle_{\text{max}} = 0.72\text{keV}$ ).

# Double Scale Length



Core Heating Properties...1D FP  $L_{f1}=1$  or  $1.5\mu\text{m}$  ( $n_e > n_c$ ) +  $L_{f2}=5\mu\text{m}$  ( $n_e < n_c$ )



In the double scale case, low energy component ( $E < 2\text{MeV}$ ) is increased, so that the energy coupling from laser to core and resultant core temperature become high compared with the single scale length cases.

$L_{f1} = 1.5\mu\text{m} + L_{f2} = 5\mu\text{m}$        $\eta_{L, \text{fuel}} = 16.7\%$  and  $\langle T_i \rangle_{\text{max}} = 0.95\text{keV}$ .

Pre-plasma scale-length dependence of core heating;

.....  $10^{20}\text{W}/\text{cm}^2$  / 750fs /  $1.06\mu\text{m}$  pulse

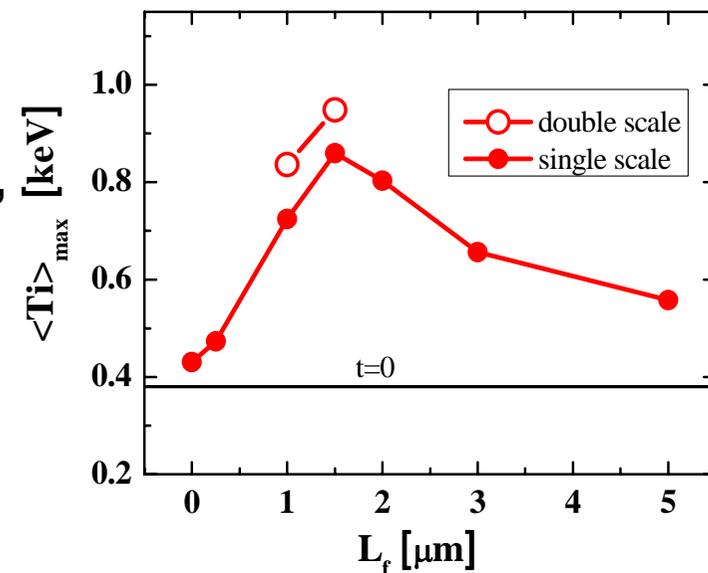
- When scale length is too short ( $L_f < 1\mu\text{m}$ ), the energy coupling from laser to fast electron is small and then the core heating efficiency is low.
- In the longer scale length region ( $L_f \gg 1\mu\text{m}$ ), with increasing  $L_f$ , the low energy component of fast electron, that mainly contributes to the core heating, becomes small, so that the heating efficiency also becomes low.
- Optimum scale length is  $1.5\mu\text{m}$

$\eta_{L \text{ fuel}} \sim 15\%$ ,  $\langle T_i \rangle$  reaches  $0.86\text{keV}$

Double scale length case;

- Energy coupling from laser to fast electron is higher than that in the single scale length case and low energy component of fast electron also increases. Core heating efficiency becomes high.

$L_{f1} = 1.5\mu\text{m}$  ( $n_e > n_c$ ) +  $L_{f2} = 5\mu\text{m}$  ( $n_e < n_c$ );  $\eta_{L \text{ fuel}} \sim 17\%$ ,  $\langle T_i \rangle_{\text{max}} = 0.95\text{keV}$



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## Longer & Higher energy Laser Case ( $>10\text{ps}$ , $>10^{20}\text{W/cm}^2$ )

- In FIREX-I, heating laser is  $10\text{kJ}/10\text{ps}$ . For ignition,  $100\text{kJ}$  (or more?) heating laser is required. What's happen? Can we Extrapolate the short & relatively low energy shot results?

## Additional Heating due to Fast Ion ?

- Accelerated ions (Au,  $Z=50$ ) does not contribute to the core heating because of their short range.
- But, how about the different ion species (C, P, D, T)?

## Geometrical effects in laser-cone interactions (e.g. *laser-cone wall interaction, laser focusing, fast electron accumulation*)

- • • • On going by T. Nakamura

## Collision processes in the cone

- • • • reduced model is under construction.

## Effects of Microscopic instabilities on fast particle transport

- In ps-order 1D PIC simulations, we observed ion heating (or acceleration) in the medium-density propagation region (between cone tip and core,  $n \sim$  a few  $10n_c$  region) due to the instability.