

## Modelling the Neutraliser Plasma of ITER Negative Ion Beams

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Plasma, Gas Heating and Neutralisation Efficiency I

• Gas heating is major contribution to the reduction in neutralisation efficiency







## Plasma, Gas Heating and Neutralisation Efficiency II

Gas heating is mainly due to collisions of gas with plasma electrons and hot neutrals

Hot neutrals formed from plasma ions accelerated across the sheath and reflected as energetic neutrals at wall

To predict gas heating we need to predict plasma conditions







The Neutraliser Model I - Source Terms

N TT

Beam evolution through the neutraliser

$$\frac{\mathrm{d}n^{\mathrm{p}}}{\mathrm{d}z} = N \Big[ n^{\mathrm{q}} \sigma_{\mathrm{qp}} + n^{\mathrm{r}} \sigma_{\mathrm{rp}} - n^{\mathrm{p}} \Big( \sigma_{\mathrm{pq}} + \sigma_{\mathrm{pr}} \Big) \Big]$$

p,q,r represent charge states -1, 0, +1

$$S_{x} = \frac{NI_{b}}{A_{b}e} \left[ f^{-}(z)\sigma_{x\bar{1}} + f^{0}(z)\sigma_{x0} + f^{+}(z)\sigma_{x1} \right]$$

Power to plasma from beam

$$P = \frac{NI_{b}\Delta z}{e} \left[ f^{-}(z)\varepsilon^{-} + f^{0}(z)\varepsilon^{0} + f^{+}(z)\varepsilon^{+} \right]$$



**Ionisation term** 



The Neutraliser Model II - Boundary Conditions

$$j_{i} - \frac{j_{e}}{4} exp\left(j_{i} - \frac{\phi j_{e}}{T_{e}^{4}}\right) exp\left(j_{e} - \frac{\phi j_{e}}{T_{e}^{4}}\right) = j_{cx} - j_{s}$$

Exact form depends upon degree of thermalisation of stripped electrons with plasma



Thermalised, 
$$E_{sB} \sim T_e$$
  $j_i - \frac{(j_e + j_s)}{4} exp\left(-\frac{\phi}{T_e}\right) = j_{cx} - j_s$ 





The Neutraliser Model - Gas Heating

J. Paméla, Rev. Sci. Instrum., 57, 1066 (1986)

$$A_{w}\alpha \frac{Nv}{4}k(T-T_{w}) = (\gamma - 1)[S_{b}(I_{b}, E_{b}) + S_{e}(n_{e}, T_{e}) + S_{i}(n_{i}, \phi, R)]$$

- $\alpha$  accommodation coefficient for gas molecules on wall
- R ion reflection coefficient
- $\boldsymbol{\gamma}$  specific heat capacity of neutraliser gas





Benchmark - Application to JET Neutraliser





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## Application to ITER Negative Ion Beams



JET



HNB 1MeV D<sup>-</sup>  $\sigma_{i\bar{1}}$  ,  $\sigma_{e\bar{1}} = 0$   $\epsilon^- = 0$ 







HNB 1MeV D<sup>-</sup> 
$$\sigma_{i\bar{1}}, \sigma_{e\bar{1}} = \sigma_{i0}, \sigma_{e0}$$
  $\epsilon^- = \epsilon^0$ 





EUROPEAN FUSION DEVELOPMENT AGREEMENT

DNB 100keV H<sup>-</sup> 
$$\sigma_{i\bar{1}}, \sigma_{e\bar{1}} = \sigma_{i0}, \sigma_{e0}$$
  $\epsilon^- = \epsilon^0$ 





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Thermalisation of Stripped Electrons I

Initial energy of stripped electrons  $E_s = \frac{m_e}{m_i} E_b$  HNB  $E_s$ =272eV DNB  $E_s$ =54.5eV

For  $E_s >> T_e$  energy loss is dominated by inelastic collisions

HNB energy loss per stripped electron 31eV,  $T_e \sim 25eV$ 

DNB energy loss per stripped electron 16.4eV,  $T_e \sim 20eV$ 

For DNB coulomb collisions will be significant and stripped electrons will be thermalised





**Thermalisation of Stripped Electrons II** 

Boundary condition 
$$j_i - \frac{(j_e + j_s)}{4} exp\left(-\frac{\phi}{T_e}\right) = j_{cx} - j_s$$



**z**(**m**)





Conclusions

- Plasma density in the neutraliser is low
  - Division into four channels reduces beam current per channel
- In HNB stripped electrons do not influence plasma
- In DNB conditions may give rise to negative plasma potential
- Gas heating is not expected to be severe in the ITER beams

