



Fishbones Activity in JET Low Density Plasmas

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Nuclear

<u>Outline</u>







Introduction

Fishbones (and sawteeth) are caused by the internal kink mode

They correspond to different solutions of the dispersion relation

Dispersion relation for the internal kink mode:

$$\delta W_{MHD} + \delta W_{HOT} - \frac{8i\Gamma[(\Lambda^{3/2} + 5)/4] [\omega(\omega - \omega_{*i})]^{1/2}}{\Lambda^{9/4} \Gamma[(\Lambda^{3/2} - 1)/4] \omega_A} = 0$$

(R. White et al. Phys. Fluids B 2, 745 (1990))

(Large aspect ratio aproximation)







Usual solutions for the internal kink dispersion relation

Unstable branches Instability Frequency

High eta_h		Fishbone branch	Precessional fishbones	$\omega \approx \omega_{D}$
Low β_h	$\gamma_I/\omega_{k}>1/2$	Kink branch	Sawteeth	
	γ_I/ω_{*i} <1/2	Ion branch	Diamagnetic fishbones	$\omega \approx \omega_{*i}$
Intermediate β_h				







New fishbones regimes in low density plasmas











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The different instabilities occur in sequence !







Objective: explain observed fishbones behaviour

Solutions of the dispersion relation depend on three parameters γ_I Ideal growth rate

- \mathcal{O}_{*i} Diamagnetic frequency
- β_h Fast particles beta

Step 1: Determine the range of parameters for which each branch of the dispersion relation is unstable.

Step 2: Determine the evolution of each parameter during the monster sawtooth.

Step 3: Determine how the evolution of these parameters affect fishbones behaviour







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Step 1: Determination of the regions of stability

To determine the regions of stability we need only to know when mode stability changes, i.e, $Im(\omega)=0$

Introducing this condition in the dispersion relation,

$$\gamma_{I} = \frac{3}{4} \left[\frac{\omega}{\langle \omega_{D} \rangle} \left(\frac{\omega}{\langle \omega_{D} \rangle} - \frac{\omega_{*i}}{\langle \omega_{D} \rangle} \right) \right]^{\frac{1}{2}} \left(\frac{\omega}{\langle \omega_{D} \rangle} \right)^{-\frac{3}{2}} \left[\frac{1}{2} + \frac{\omega}{\langle \omega_{D} \rangle} + \left(\frac{\omega}{\langle \omega_{D} \rangle} \right)^{\frac{3}{2}} \operatorname{Re} Z \left[\left(\frac{\omega}{\langle \omega_{D} \rangle} \right)^{\frac{1}{2}} \right] \right]$$

$$\frac{\beta_{h}}{4} = \frac{3}{4} \frac{\varepsilon \omega_{A}}{\pi^{1/2} \langle \omega_{D} \rangle} \left[\frac{\omega}{\langle \omega_{D} \rangle} \left(\frac{\omega}{\langle \omega_{D} \rangle} - \frac{\omega_{*i}}{\langle \omega_{D} \rangle} \right) \right]^{\frac{1}{2}} e^{\omega/\omega_{D}} \left(\frac{\omega}{\langle \omega_{D} \rangle} \right)^{-\frac{5}{2}}$$
Two solutions
$$\gamma_{I} < \gamma_{M}$$
(ideal limit, ICRH population)
And ω_{*i} takes the role of a parameter











Step 2: Temporal evolution of the parameters



Decreases during a burst Increases between bursts









Step 3: Changes due to the evolution of the parameters









 β_h

During the monster sawtooth, the region where the high and low frequency modes coalesce is reached.







Solutions of the dispersion relation





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Resonant transference of energy between between an internal kink mode (top hat structure) and an ICRH driven fast ions population as function of the frequency



















Summary and Conclusions

In low density plasmas the diamagnetic frequency has a significant increase:









Summary and Conclusions



Fast ions temperature between 1 MeV and 1.5 MeV maximize the resonant exchange of energy for frequencies 50-80 kHz.



Numerical simulations predict a rapid increase in the mode frequency as the fast ions temperature increase.



This may be related with a change in the type of orbits that have stronger interaction with the mode.

