

Simulations on the Nonlinear Mode Coupling in Multiple-scale Drift-type Turbulence with Coherent Flow Structures

Jiquan Li^{1,2},

K. Uzawa², Z. Lin³, Y. Kishimoto², N. Miyato⁴, T. Matsumoto⁴, J.Q. Dong¹

- 1) Southwestern Institute of Physics, Chengdu, China
- 2) Graduate School of Energy Science, Kyoto University, Uji, Japan
- 3) Department of Physics and Astronomy, University of California, Irvine, USA
- 4) Fusion Research and Development Directorate, JAEA, Japan

e-mail: lijq@energy.kyoto-u.ac.jp

Acknowledgements: L. Chen, P. Diamond, H. Sanuki, J. Anderson, M. Azumi, Y. Liu, M. Kikuchi, H. Ninomiya

Motivation & Main Results

» Multi-scale turbulence and secondary/tertiary coherent flow structures coexist in plasmas (MHD; ITG; TEM; ETG;..... Mean/zonal flows; streamers; KH, GKH modes;)

» Flow characteristics and interaction mechanisms





Wave-type flow with low/zero frequency: nonlinear mode coupling ?



Li/Kishimoto PoP 2004

» Motivation:

Secondary and/or tertiary structures as a wave-type flow interact with turbulence through a primary mechanism: nonlinear mode coupling

» Main results

► Back-action of zonal flows/long wavelength modes on turbulence deforms the power-law spectrum into an exponential-law scaling.

Secondary streamer-like long wavelength modes can saturate ETG turbulence, suggesting the possibility of a low ETG fluctuation level and electron transport. **Spectral characteristics of turbulence due to back-action of flow structures**

Schematic picture of nonlinear mode coupling

» Ubiquitous nonlinear process: basic four-wave & three-wave interaction



Simulations in 2D forced HM model vs GTC

» Forced HM turbulence modeling based on 2D ETG (Free energy system)

$$(1 - \nabla_{\perp}^{2})\frac{\partial \Phi}{\partial t} = \frac{\partial \Phi}{\partial y} + [\Phi, \nabla_{\perp}^{2}\Phi] + \gamma_{drive}^{ETG}(\vec{k})\Phi + \gamma_{damp}^{ETG}(\vec{k})\Phi$$

» Parameters:



Identifying mode coupling due to zonal flows

» k_x spectral relation between ETG fluctuations and zonal flows due to nonlinear mode coupling



Sen et al. PoP 2006

4-wave coupling for zonal flow generation (modulation instability) also acts back on turbulence to produce high k_x fluctuation which is damped. Larger k_q and wider zonal flow spectrum can make larger k_x spectral hump. Interaction becomes nonlocal.

Spectral transition due to KH/GKH mode

2D forced HM modeling ETG simulation



Nonlinear mode coupling between KH mode with $k_y \approx 0.1$ interacts with ETG $(k_y \approx 0.4)$ deform power-law spectrum (t<1500) into exponential scaling(t>1500) **Gyrokinetic ETG simulation by performing GTC code**



Similar to fluid simulation, long wavelength GKH is excited before ETG simulation; Exponential-law spectrum is also observed after ETG saturation.

Spectral characteristics in 3D slab ETG / ITG

» Simulations with and no zonal flows for comparison



» Without zonal flows: k_x or k_y power-law spectrum $\phi^2(k_x) \sim k_x^{-6.3}$

same as theoretical estimate $\alpha = 6.29$ [Ottinger & Carati, PRE(1993)]

» With zonal flows: exponential-law $\phi^2(k_x) \sim e^{-3.0k_x}$, almost no change for k_y spectrum.

Experimental exponential-law density spectra in Tore Supra, Hennequin, et al. TTF06

Back-action of zonal flows through nonlinear mode coupling may provide a drive force to deform power-law scaling into an exponential-law dependence **Role of streamer-like long wavelength structures in ETG saturation**

Electron transport vs large-scale flow in ETG

» Several ETG simulations show streamer structures, but transport is different.
» Streamer structure is expected to enhance electron transport. However, what is its probable role in ETG saturation?

- » Assuming wave-type streamer-like long-wavelength mode imposed in ETG (ignoring k_x) $\phi_s = \phi_{s0}(t) \cos(k_s y)$ $k_s <<1$
- » Coupled equs. due to streamer-like long-wavelength mode

$$\begin{cases} \partial_{t} (1 - \nabla_{\perp}^{2}) \phi^{k_{y}} = (1 + K \nabla_{\perp}^{2}) \partial_{y} \phi^{k_{y}} - \frac{i}{2} k_{s} \phi_{s} \partial_{x} \nabla_{\perp}^{2} (\phi^{k_{y}+k_{s}} - \phi^{k_{y}-k_{s}}) + \nabla_{\parallel} v_{\parallel}^{k_{y}} - \mu_{\perp} \nabla_{\perp}^{4} \phi^{k_{y}} \\ \partial_{t} v_{\parallel}^{k_{y}} = \nabla_{\parallel} (\phi^{k_{y}} - p_{e}^{k_{y}}) + \frac{i}{2} k_{s} \phi_{s} \partial_{x} (v_{\parallel}^{k_{y}+k_{s}} - v_{\parallel}^{k_{y}-k_{s}}) + \eta_{\perp} \nabla_{\perp}^{2} v_{\parallel}^{k_{y}} \\ \partial_{t} p_{e}^{k_{y}} = -K \partial_{y} \phi^{k_{y}} - \frac{5}{3} \nabla_{\parallel} v_{\parallel}^{k_{y}} + \frac{i}{2} k_{s} \phi_{s} \partial_{x} (p_{e}^{k_{y}+k_{s}} - p_{e}^{k_{y}-k_{s}}) + \sqrt{\frac{32}{9\pi}} |k_{\parallel}| (p_{e}^{k_{y}} + \phi^{k_{y}}) + \chi_{\perp} \nabla_{\perp}^{2} p_{e}^{k_{y}} \end{cases}$$





 k_y -mode coupling in slab corresponds to toroidal mode coupling in tokamak since slab is the local expansion near given q_0 surface.

Z. Lin / L. Chen et al. IAEA 2004, PoP2005, PPCF 2005

3D ETG simulation vs long-wavelength mode





Excitation of long wavelength mode in ETG

+c.c.

» Hasegawa-Mima turbulence model

 $(1 - \nabla_{\perp}^{2}) \frac{\partial}{\partial t} \phi = \frac{\partial}{\partial y} \phi + [\phi, \nabla_{\perp}^{2} \phi]$

5-wave modulation: 2 pumps, 2 sidebands and flow seed

$$\phi_{p1,2} = \phi_{1,2} e^{i\vec{k}_{1,2}\cdot\vec{x} - i\omega_{1,2}t} + c \cdot c \cdot \phi_q = \widetilde{\phi}_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = \widetilde{\phi}_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = \widetilde{\phi}_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = \widetilde{\phi}_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = \widetilde{\phi}_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = \widetilde{\phi}_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = \widetilde{\phi}_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = \widetilde{\phi}_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = \widetilde{\phi}_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} + c \cdot c \cdot \phi_q = c \cdot \phi_q e^{i\vec{k}_q\cdot\vec{x} - i\omega_q t} +$$



» Dispersion relation of modulation instability

$$F(\Omega_q; \tilde{\phi}_1; \tilde{\phi}_2; \vec{k}_1; \vec{k}_2) = \mathbf{0}$$

- Low amplitude threshold for modulation instability
- Weak instability with peaks at k_x=0.8 and 0.35
- ► Wave number matches



ETG saturation due to streamer-like mode

» Streamer-like mode as a beat wave in pure linear ETG fluctuations $\phi_s = 0.1 | \phi(t)_{k_y=0.5} \phi(t)_{k_y=0.6} | \cos(0.1y)$



Time-dependent long wavelength streamerlike mode interact with ETG modes to produce k_y -mode coupling so that ETG is saturated at lower level. **» ETG saturation level vs intensity of streamer-like mode**

 $\phi_s = f^* | \phi(t)_{k_y=0.5} \phi(t)_{k_y=0.6} | \cos(0.1y)$



ETG saturation amplitude inversely proportional to intensity of streamerlike mode

Comparison of ETG saturation features

» Same physical parameters and numerical settings: $\eta_e = 6, \hat{s} = 1.4, \mu_{\perp} = \eta_{\perp} = \chi_{\perp} = 0.5; L_x = 50 \rho_e, L_y = 20 \pi \rho_e, L_z = 2\pi L_n$



2D modeling simulation can well reproduce the main features of 3D ETG turbulence saturation (amplitude, time evolution) Secondary long wavelength streamer-like flow saturate ETG turbulence through nonlinear mode coupling!

Summary and conclusion

Nonlinear mode coupling is testified and emphasized as a primary interaction mechanism in multiple-scale drift turbulence with coherent flows, which are characterized by wave-type structures. 2D modeling analysis and 3D gyrofluid slab ETG simulations well produce spectral features observed in gyrokinetic particle ETG simulation by using GTC code.

Main results:

» Through nonlinear mode coupling , zonal flow or long wavelength KH (GKH) can deform the conventional power-law k_x or k_y spectra into an exponential-law scaling.

» Secondary long-wavelength streamer-like structure can saturate slab ETG turbulence at lower level through producing a k_y -mode coupling, suggesting the possibility of low electron transport in ETG turbulence.

In addition, the effect of wave-type mean flows on the zonal flow generation in drift wave turbulence has also been investigated, the results will be presented in our poster: TH/2-3

Thank you!